STRANGE HADRONS: Bridge WED QUARKS And SLARS

• MOtivation.

• HADRONS IN COLDLNUCLEI • EXOTIC HYPERNUCLEI

• Double hypernuclei

- ANTIBARYONS in nuclei

conclusion





Pillars of Hadron Physics





What is the heart of a Neutronstar?



- ...even more speculations: supersymmetric baryons
 S. Balberg *et al.*, The Astrophysical Journal, 548:L179–L182 (2001)b
- present data on neutron star masses do not (yet) exclude exotic cores
- simultaneous treatment of all possible ingredients (K,Y,q...) missing

Moderne Version of B-B Interaction

 Baryons and their mutual interactions is a complex, quantumfieldtheoretical, non-perturbative many-body problem



exotic hypernuclei

Baryon-baryon interaction





► Isospin dependence of Y-N and Y-Y interaction? ⇒ Information on hyperons in neutron rich matter/nuclei needed

International Hypernuclear Network

PANDA at FAIR Anti-proton beam Double Λ-hypernuclei

 γ -ray spectroscopy

MAMI C

- Electro-production
- Single Λ -hypernuclei
- Λ-wavefunction

JLab

- Electro-production
- Single Λ -hypernuclei
- Λ -wavefunction

FINUDA at DAΦNE

- e⁺e⁻ collider
- Stopped-K⁻ reaction
- Single Λ -hypernuclei
- γ-ray spectroscopy

HypHI at FAIR Heavy ion beam Single Λ-hypernuclei At extreme isospins Magnetic moments

J-PARC

- Intense K⁻ beam
- Single and double Λ -hypernuclei
- γ -ray spectroscopy for single Λ

Relativistic Hypernuclei





The HYPHI Project T. Saito





double hypernuclei

Production of $\Lambda\Lambda$ -Hypernuclei



- simultaneous implantation of two Λ is not feasible
- ▶ reaction with lowest Q-value: $\Xi^{-}p \rightarrow \Lambda\Lambda$: 26MeV
- ► direct implantation of a Ξ⁻ via a two-body reaction difficult because of large momentum transfer



spectroscopie BSOJEE BIDE IM ONES INTENDE MIENTON CONTRACTOR Products

KEK-E373: the NAGARA event



- ▶ H. Takahashi *et al.*, PRL 87, 212502-1 (2001)
 - hybrid emulsion technique
 - cleanest event so far (also theoretically)

$$\Xi^{-} + {}^{12}C \rightarrow {}^{4}He + t + {}^{6}_{\Lambda\Lambda}He$$
$${}^{6}_{\Lambda\Lambda}He \rightarrow {}^{5}_{\Lambda}He + p + \pi^{-}$$
$$\Rightarrow \Delta B_{\Lambda\Lambda} = +1.01 \pm 0.2^{+0.18}_{-0.11}MeV$$

- no electronic detector can compete with emulsions
- structure with high resolution ⇒
 γ-spectroscopy with Ge detectors



The Discovery of the anti-Xi





PHYSICAL REVIEW LETTERS

MARCH 15, 1962





VOLUME 8, NUMBER 6

OBSERVATION OF PRODUCTION OF A $\Xi^+ \Xi^+$ PAIR*

H. N. Brown, B. B. Culwick, W. B. Fowler, M. Gailloud,[†] T. E. Kalogeropoulos, J. K. Kopp, R. M. Lea, R. I. Louttit, T. W. Morris, R. P. Shutt, A. M. Thorndike, and M. S. Webster Brookhaven National Laboratory, Upton, New York

and

C. Baltay, E. C. Fowler, J. Sandweiss,[‡] J. R. Sanford, and H. D. Taft Yale University, New Haven, Connecticut (Received February 19, 1962)

VOLUME 8, NUMBER 6

PHYSICAL REVIEW LETTERS

MARCH 15, 1962

EXAMPLE OF ANTICASCADE ($\overline{\Xi}^+$) PARTICLE PRODUCTION IN $\overline{\rho}$ - ρ INTERACTIONS AT 3.0 GeV/c

CERN, Geneva, Switzerland * Laboratoire de Physique, Ecole Polytechnique, Paris, France

and

Centre d' Etudes Nucléaires, Département Saturne, Saclay, France

(Received February 19, 1962)

An experiment is in progress at the CERN proton synchrotron to study the interactions of fast antiprotons with protons. A high-energy separated beam¹ has been installed and optimized to provide, in the first instance, a high-purity beam of 3.0-Gev/c antiprotons. The interactions are being produced and observed in the Saclay 81-cm hydrogen bubble chamber.²

In the methodical scanning of the first ten thousand photographs (with an average of seven antiprotons per photograph) an event has been found showing the production of an anticascade particle $(\overline{\Xi}^+)$. The object of this Letter is to present the data and the analysis leading to this conclusion.

One of the three views of the event is reproduced in Fig. 1. Briefly, the event is as follows: After travelling 20 cm in the chamber, a beam particle

interacts at point A, producing two charged particles. The positive particle decays at point B (distant 6 cm from A) and the negative at point D (4 cm from A). Both decay secondaries are light particles, as we will see. At C-about 20 cm downstream from B-there appears a V^0 , which will be identified later as the decay of a $\overline{\Lambda}^{o}$ particle. Near point B another two-prong interaction can be seen at point I: Stereoscopic reconstruction shows that there is no direct link between this interaction and the $\overline{\Lambda}^{0}$ decay.

The event can be analyzed in several ways. We have chosen to proceed in two steps: We first analyze the event connected with the positive particle from apex A, and then with the improved knowledge thus derived we analyze the complete interaction at the same apex.

Production of Ξ^-



► Ξ^- conversion in 2 Λ : $\Xi^- + p \rightarrow \Lambda + \Lambda + 28.5 \text{MeV}$



antiproton storage ring HESR

 $p + \overline{p} \rightarrow \Xi^{-} + \overline{\Xi}^{+}$ $p + \overline{p} \rightarrow \Omega^{-} + \overline{\Omega}^{+}$

• few times 10^5 stopped Ξ per day $\Rightarrow \gamma$ -spectroscopy feasible



PANDA setup





Can we do more?



- triple hypernuclei via $p\overline{p} \rightarrow \Omega\overline{\Omega}$ $\Omega pn \rightarrow \Lambda\Lambda\Lambda + 203 MeV$?
 - Iower cross section
 - large momenta \Rightarrow lower stopping probability
 - large Q-value \Rightarrow low probability for triple Λ nuclei
 - γ -spectroscopy most likely not practical at the beginning

• Λ_c hypernuclei

- production via primary + secondary target not possible because of short lifetime of $\tau_{\Lambda c}$ =0.2ps which exceeds stopping time
- direkt production via pp $\rightarrow \Lambda_c \Lambda_c bar$ or $\pi^- p \rightarrow \Lambda_c D^-$ difficult because of high momenta involved (very low sticking probability)
- does a two-step process within one nucleus work?

$$\overline{p} + p \rightarrow D^{+} + D^{-}$$
 detected

$$D^{+} + p \rightarrow \Lambda_{c}^{+} + \pi^{+}$$
captured in the
nucleus A-2

- \blacktriangleright determination of the $\Lambda_{\rm c}$ hypernucleus mass via missing mass
 - \triangleright needs good knowledge of beam momentum (10⁻⁴)
 - ▷ excellent momentum resolution for π^+ and D⁻ (resp. decay products)
- expected rate ~0.01 day⁻¹ (??? rescattering \rightarrow 1day⁻¹???)

antibaryons in nuclei

G-Parity and NN Potential

GUTENBERG MANZESTÄT

- strong interaction conserves isospin and C-parity
- G=charge conjugation + 180° rotation around 2nd axis in isospin
 - Lee und Yang 1956, L. Michel 1952 "Isoparity"
 - G-parity of particle-antiparticle multipletts

 $G | \mathbf{f} \overline{\mathbf{f}} \rangle = (-1)^{I} C | \mathbf{f} \overline{\mathbf{f}} \rangle = (-1)^{I+L+S} | \mathbf{f} \overline{\mathbf{f}} \rangle$ $G | \pi^{\pm 0} \rangle = (-1)^{1} C | \pi^{\pm 0} \rangle = - | \pi^{\pm 0} \rangle$ $G | \rho \rangle = (-1)^{1} C | \rho \rangle = + | \rho \rangle$ $G | \omega \rangle = (-1)^{0} C | \omega \rangle = - | \omega \rangle$ $G | \sigma \rangle = (-1)^{0} C | \sigma \rangle = + | \sigma \rangle$

- Hans-Peter Dürr and Edward Teller, Phys. Rev. 101, 494 (1956)
 - sign change in coupling constant

$$V(NN)(r) = \sum_{M} V_{M}(r) \rightarrow V(N\overline{N})(r) = \sum_{M} G_{M}V_{M}(r)$$



Elastic Antiproton-Nucleus Scattering

Elastic Scattering of Antiprotons from Complex Nuclei*

Gerson Goldhabert and Jack Sandweiss‡

Physics Department and Radiation Laboratory, University of California, Berkeley, California (Received May 5, 1958) pantiwith $\geq 2^{\circ}$.)

TABLE III. Comparison of experimental data for elastic antiproton-nucleus scattering of energy $T_{\bar{p}}=80$ to 200 Mev with Glassgold's calculations at $T_{\bar{p}}=140$ Mev. (Projected angle $\geq 2^{\circ}$.)

Angular interval (degrees)	Experimental $(T_{\overline{p}} = 80 \text{ to}$ 200 Mev)	Number of events Calculated for V = -15 Mev W = -50 Mev	Calculated for V = -528 Mey W = -50 Mey
2-6 6-12 12-24 24-180	54 20 5 1	$56 \\ 17.1 \\ 4.3 \\ 1.4$	71 24 10 9.5
2–180	80	78.8	114.5

Antiprotonproduction in HI Collisions

see e.g.

A. Sibirtsev, W. Cassing et al., Nucl. Phys. A 632, 131 (1998)

C. Spieles et al., Phys. Rev. C 53, 2011-2013 (1996)



 $p + {}^{12}C \longrightarrow \overline{p} + X$ at 4.0 GeV

Can we measure the potential for \overline{Y} ?

- ▶ $p + \overline{p} \rightarrow \Lambda + \overline{\Lambda}$ close to threshold within complex nuclei
- Question: is the momentum of the Λ and anti-Λ equal?
- ► If yes, Λ and anti-Λ that leave the nucleus will have different asymptotic momenta
 - the momentum difference is sensitive to the potential difference



- experimental complications
 - Fermi motion
 - leading effect
 - exclusiveness

 \Rightarrow need to look at average transverse momentum close to threshold of coincident $\Lambda\bar{\Lambda}$ pairs

Pair production in Nuclei: $\overline{p}^{12}C \rightarrow \overline{\Lambda}\Lambda X^{\text{manual}}$

- Simulations: Antiproton momentum: 1.66 GeV
 - ▶ ∧ potential: -28MeV
 - Fermi motion (1s_{1/2} and 1p_{3/2} single-particle wf)
 - angular distribution (leading effect)
 - absorption (still crude)



A Closer Look...



Λ



Outlook



the (exclusive) production of B-B pairs in nuclei by antiproton beams may offer the possibility to study the behaviour of antibaryons in nuclei



Open tasks

- momentum dependent absorption
- momentum dependent rescattering
- formation time
- …and probably many more

conclusion

Conclusion

- Antiproton collisions with nuclei offer many opprtunities to study strange baryons in cold nuclei
 - baryon-baryon interaction
 - properties of (anti-)baryons in nuclei
 - spectroscopy of baryonic atoms
 - weak decay
 - ► ...
- These studies are made possible by a unique combination of experimental facilities at FAIR
 - ► hypernuclei spectrometer ⊕ secondary relativistic HI beams
 - γ -spectroscopy with Ge detectors \oplus antiproton beams

Baryon stars

10[°]

10⁻¹

10⁻²

10⁻³

10⁻⁴

Baryon fraction



N. K. Glendenning, Phys. Rev. C 64, 025801 (2001)

Density (fm^{-3})

- beyond $2\rho_0$ hyperons may play a significant role in neutron stars
- in the core hyperons may even be more abundant than neutrons
- needed: BB interaction at high density = at small distances

Supersymmetric Particles



- ▶ The Astrophysical Journal, 548:L179–L182, 2001
- Shmuel Balberg, Glennys R. Farrar, and Tsvi Piran



Baryon-Baryon "Potential"

■ Baryon-Stars ⇔ BB interaction at high density = at small distances



Nijmegen OBE Potential fitted to ΛN Daten + SU(3) Symmetrie Th.A. Rijken, V.G.J. Stoks, Y. Yamamoto, Phys. Rev. C 59, 21 (1999)

Observable Consequence

Main consequence of hyperons and other exotica in neutron stars: softer EOS \Rightarrow lower mass and smaller central density

U1700-37 I

ela X-1 (b,c) yg X-2 (d)

X-ray

binaries

double

binaries

neutron star

white dwarf-

neutron star

2.5

3.0

binaries

2.0

James M. Lattimer and Madappa Prakash Phys. Rev. Lett. 94, 111101 (2005)



present data on neutron star masses do not exclude exotic cores

simultaneous treatment of all possible ingredients (K,Y,g...) missing

Birth, life and death of a hypernucleus



Exp. Approaches to Y-N interactions



▷ JPARC: ~1000 events/day

- hyperon-hyperon final state interaction
 - feasible but difficult to interpret
- hyperons bound in nuclei
- hyperon-antihyperon pair production in nuclei

Ξ^- scattering



► Ahn et al.



Summary of E864



 Limits compatible with coalescence taking low biding energy and large size of hypernculei into acount

Magnetic moment of Λ in nuclei

- GUTENBERG MANZ
- Baryons do not "melt" in nuclei: quark effects are small
- EMC-effect: Whether there is any change in nucleon properties in nuclei remains controversial.
 - If mass and size of a baryons changes inside nuclei, also it's magnetic moment might change
 - If so, why? Meson current, $\Lambda \Sigma$ mixing, partial deconfinement...?



Traditional View of the N-N Interaction

- Experimental observation
 - short range (r<0.5fm)
 - intermediate (r≈1fm)
 - long range (r>1.5fm)
- Boson exchange model
 - Yukawa (1935)
 - Klein-Gordon equation

$$\left(\partial^2 + m^2\right)\varphi(x) = g\overline{\psi}\psi$$

repulsion strong attraction attraction



r [fm]

after R. Machleidt, Adv. Nucl. Phys. 19, 189 (1989)

Weak Decay of A Hypernuclei





Weak baryon-baryon interaction (1) GUTENEERS

 non-mesonic weak decay of hypernuclei explore the baryon-baryon weak interaction





- only parity violating part of weak interaction
- parity-conserving part masked by strong interaction

- parity violating and parityconserving part of weak, strangeness changing, interaction
- q~400 MeV/c
 ⇒ probes short distances

Weak decay in double Hypernuclei



- only parity violating part of weak interaction
- parity-conserving part masked by strong interaction
- parity violating and parity-conserving part of weak, strangeness changing interaction
- meson vs. direct quark process
- Interesting theoretical developments:
 - Effective Field Theories in S=-1 sector

A Parreno C Bennhold and B R Holstein nucl-th/0308074 & 0308056 weak decay studies need the detection of the decay pion or nucleon

S.R. Beane et al., nucl-th/0311027

Rate



- with increasing antiproton momentum decreasing minimum D⁺ momentum
- smaller phase space at larger antiproton momenta



▶ production rate ~0.01 day⁻¹....*hopeless!*?

Is there still a happy end possible?

• relevant for Λ_c production are D⁺ momenta < 1-2 GeV/c



P(D),P(D) [GeV/c]

Chiral Symmetry (John Cramer)

- **Question 1:** The up and down "current" quarks have masses of 5 to 10 MeV. The π^- (a down + anti-up combination) has a mass of ~140 MeV. Where does the observed mass come from?
- Answer 1: The quarks are more massive in vacuum due to "dressing". Also the pair is tightly bound by the color force into a particle so small that quantum-uncertainty *zitterbewegung* gives both quarks large average momenta. Part of the π^- mass comes from the *kinetic energy* of the constituent quarks.

Question 2: What happens when a pion is placed in a hot, dense medium? Answer 2: Two things happen:

- 1. The binding is reduced and the pion system expands because of external color forces, reducing the *zitterbewegung* and the pion mass.
- 2. The quarks that were "dressed" in vacuum become "undressed" in medium, causing up, down, and strange quarks to become more similar and closer to massless particles, an effect called "chiral symmetry restoration". In many theoretical scenarios, chiral symmetry restoration and the quark-gluon plasma phase usually go together.
- Question 3: How can a pion regain its mass when it goes from medium to vacuum?
- Answer 3: It must do work against an average attractive force, losing kinetic energy while gaining mass. In effect, it must climb out of a potential well that may be 140 MeV deep.





medium

vacuum