Exploring the Potential of Antibaryons in Nuclei with Antiprotons

Josef Pochodzalla

motivation
 proposed method
 PS185 data
 outlook

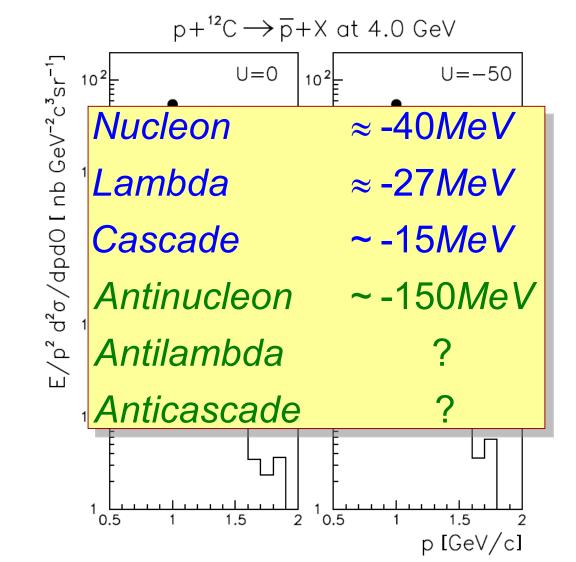
J.P., Physics Letters **B** 669 (2008) 306–310 J.P., Hyperfine Interactions, Springer, ISSN0304-3843 (Print) 1572-9540 (Online) J.P and Stephan Pomp, SENDAI08

Antiprotonproduction in HI Collisions GUTENBERG

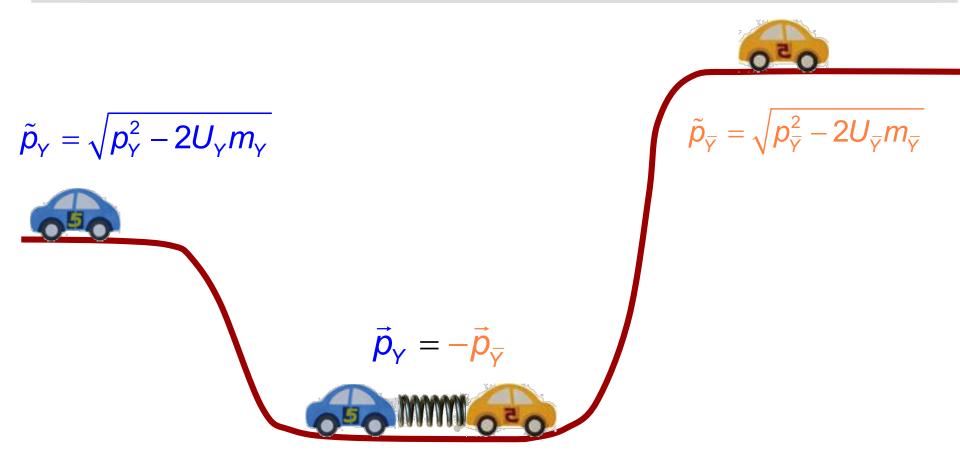
see e.g.

A. Sibirtsev, W. Cassing et al., Nucl. Phys. A 632, 131 (1998)

C. Spieles et al., Phys. Rev. C 53, 2011-2013 (1996)



How to measure a potential (difference)

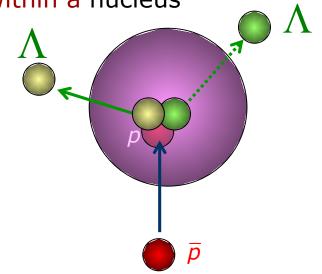


• If $m_Y \approx m_{\overline{Y}} \approx m$ and $U_Y \approx U_{\overline{Y}} \approx U \Rightarrow$

$$\alpha = \frac{\tilde{p}_{Y} - \tilde{p}_{\bar{Y}}}{\tilde{p}_{Y} + \tilde{p}_{\bar{Y}}} = \frac{\sqrt{p_{0}^{2} - 2m_{Y}U_{Y}} - \sqrt{p_{0}^{2} - 2m_{\bar{Y}}U_{\bar{Y}}}}{\sqrt{p_{0}^{2} - 2m_{Y}U_{Y}} + \sqrt{p_{0}^{2} - 2m_{\bar{Y}}U_{\bar{Y}}}} \approx \frac{U_{\bar{Y}} - U_{Y}}{4\left(\frac{p_{0}^{2}}{2m} - U\right)} \approx \frac{U_{\bar{Y}} - U_{Y}}{4E_{kin}}$$

Can we measure the potential for Y?

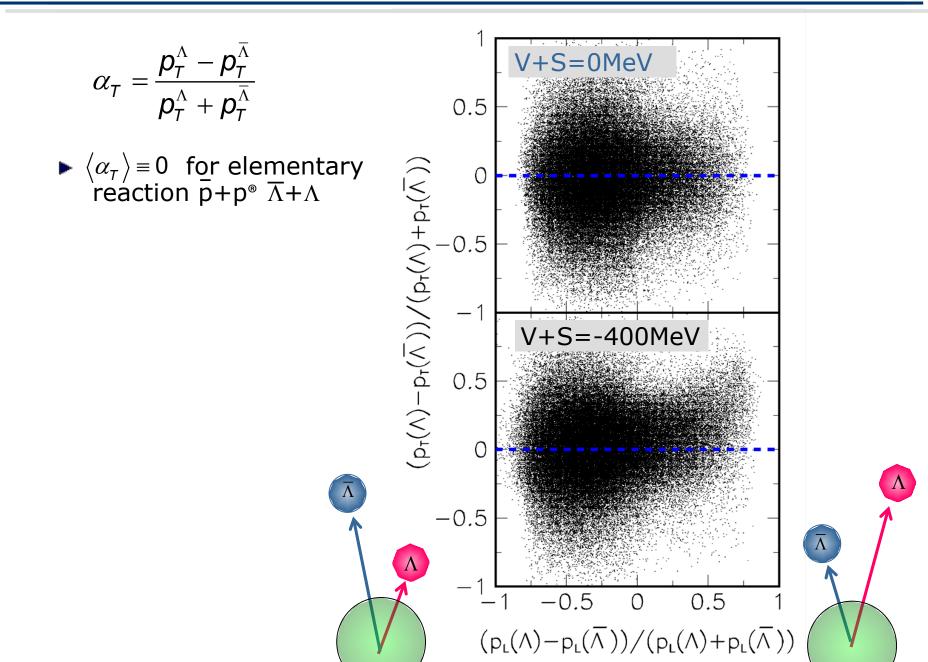
- antiprotons are optimal for the production of mass without large momenta
- consider $p + \overline{p} = Y + \overline{Y}$ close to threshold within a nucleus
- A and A that leave the nucleus will have different asymptotic momenta depending on the respective potential
- experimental complications
 - Fermi motion of struck proton
 - Non-isotropic production
 - Density distribution U(ρ)
 - Exclusiveness



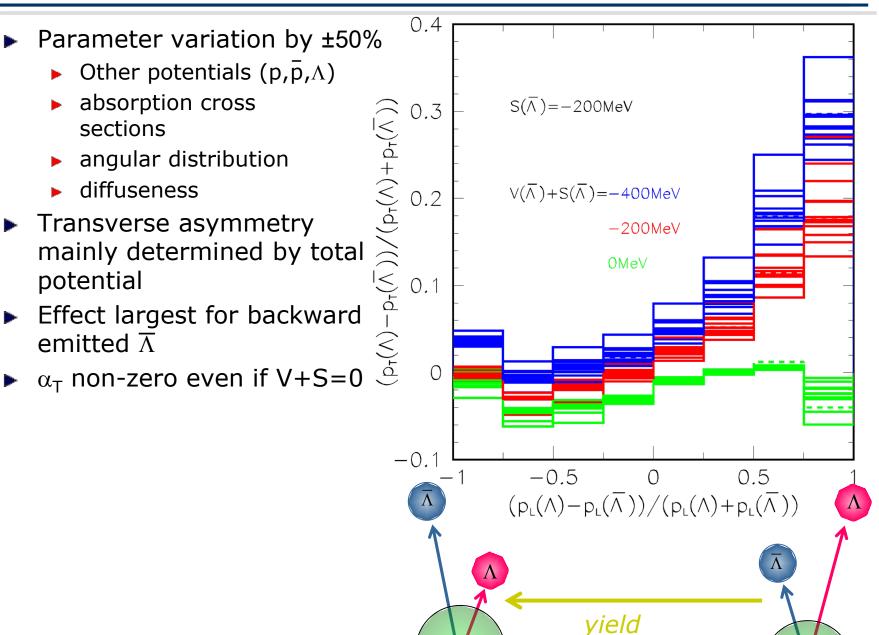
 \Rightarrow need to look at average *transverse* momentum close to threshold of *coincident* $Y\overline{Y}$ pairs

Transverse Momentum Correlations



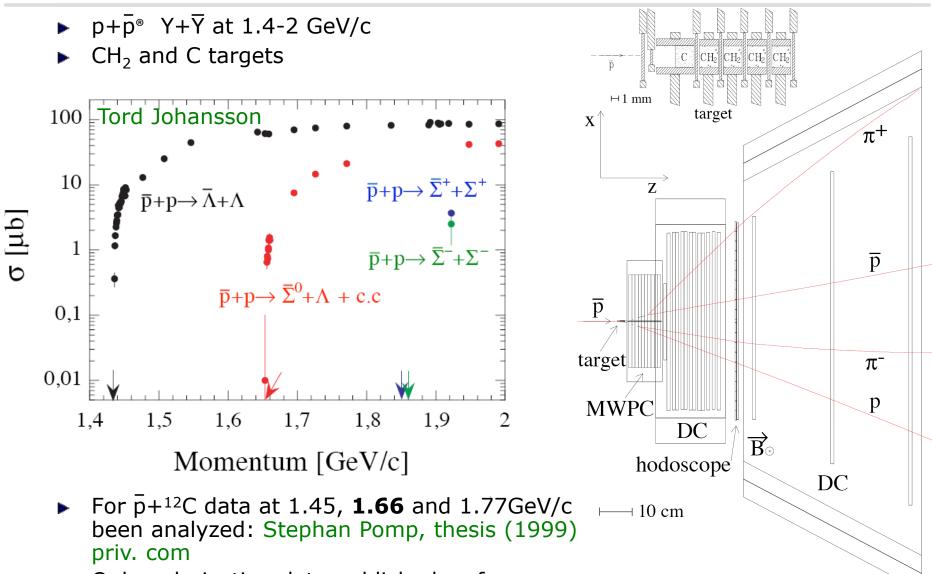


Parameter Scan



PS185 at LEAR

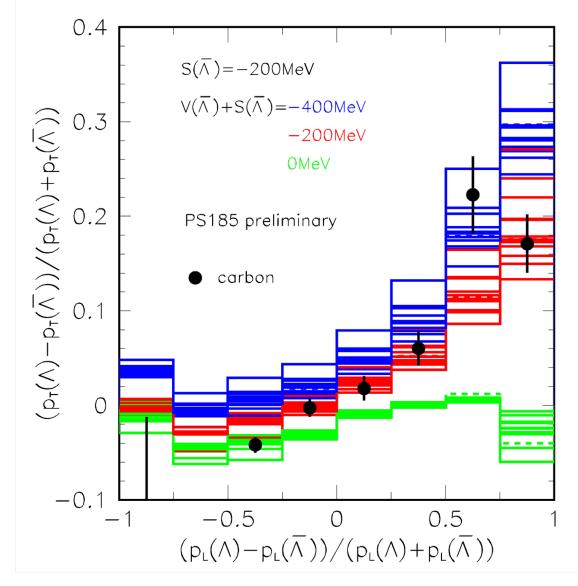




Only polarization data published so far

The good news...

Calculations are probably not incompatible with available data

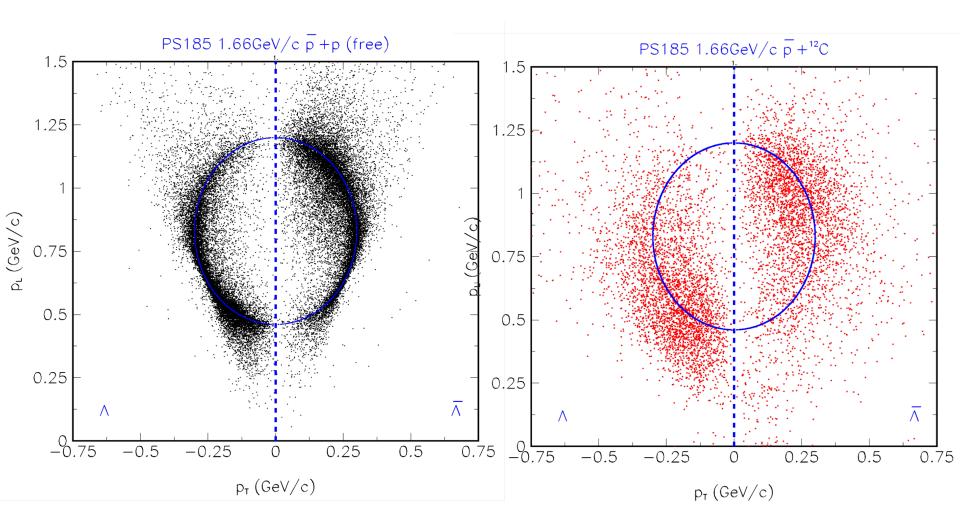




PS185: the problem

GUTENBERG MAUNVERSITÄT

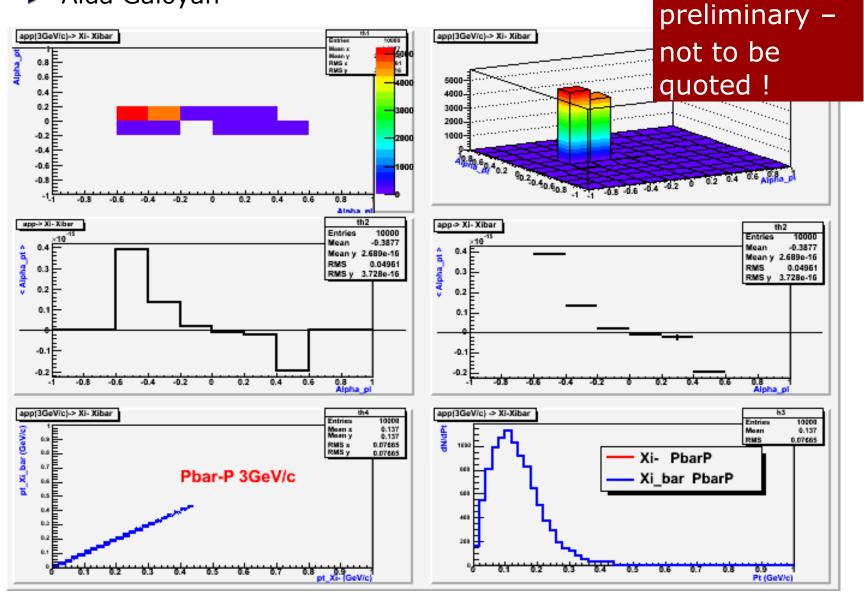
- Elementary reaction not completely understood
- But: Measurements are possible !!!



UrQMD Calculations $\Xi - \overline{\Xi}$



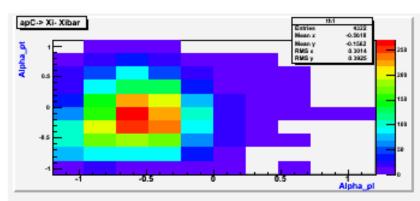
Aida Galoyan

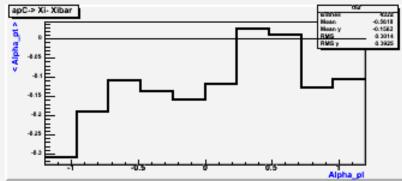


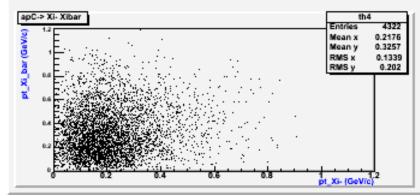
UrQMD Calculations $\Xi - \overline{\Xi}$

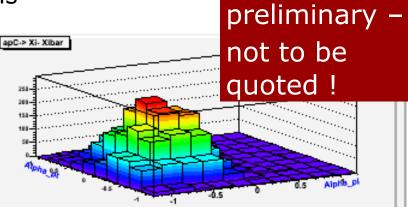


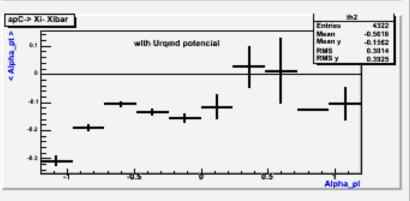
Aida Galoyan: UrQMD potentials

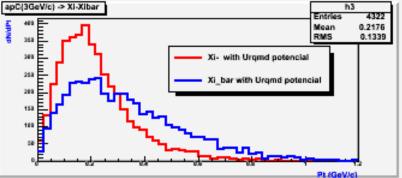












UrQMD Calculations $\Xi - \overline{\Xi}$



Inha

Pt (GeV/c)

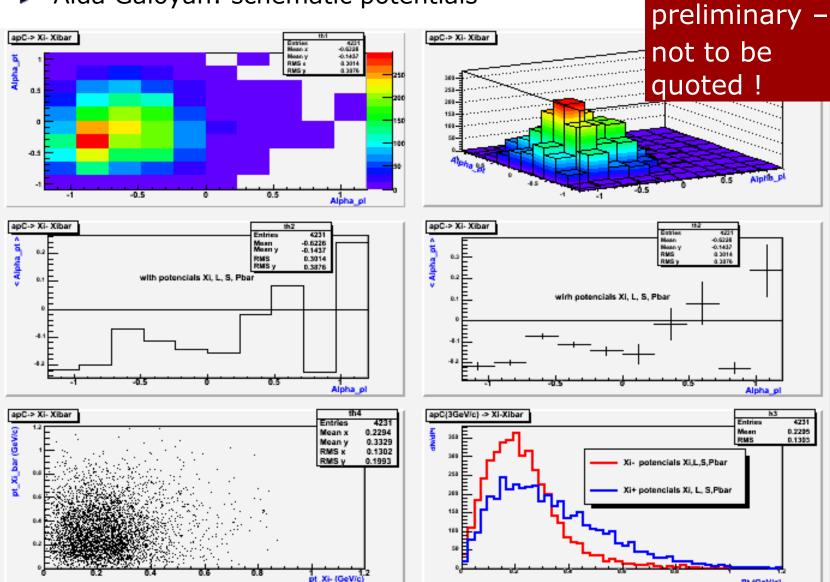
h3

4231

0.2295

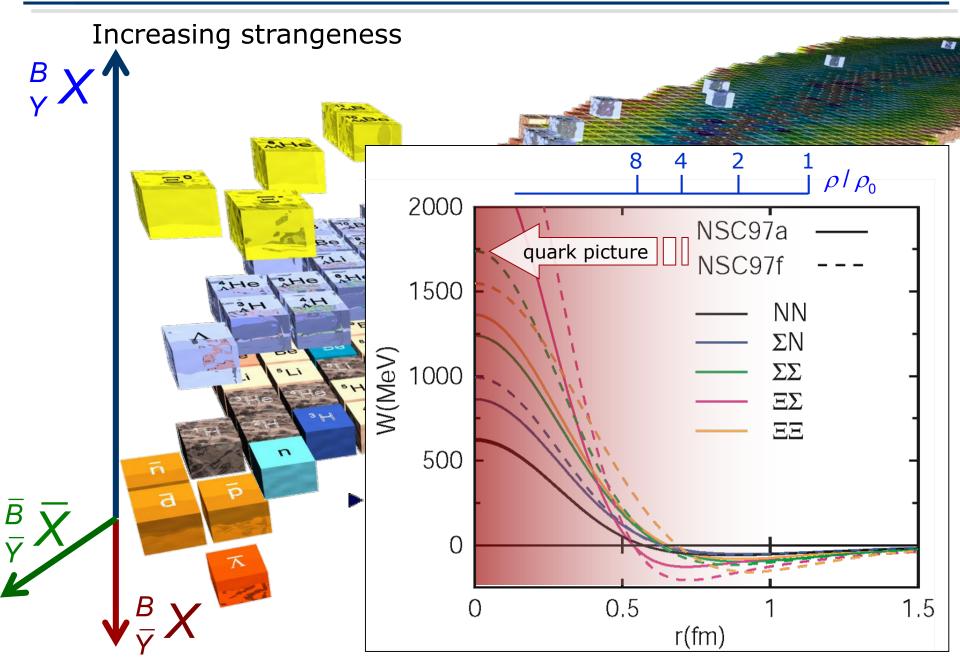
0.1303





The present nuclear chart





G-Parity and NN Potential

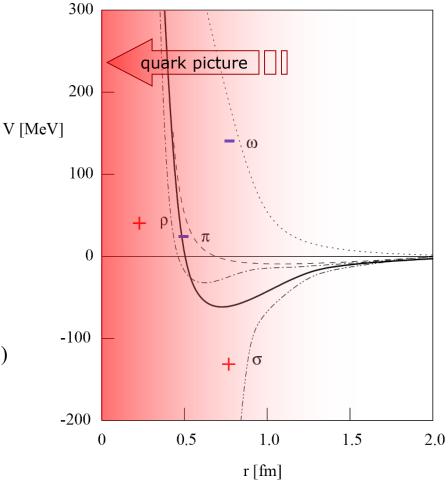
GUTENBERG MANZESTÄT

- strong interaction conserves isospin and C-parity
- G=charge conjugation + 180° rotation around 2nd axis in isospin
 - Lee und Yang 1956, L. Michel 1952 "Isoparity"
 - G-parity of particle-antiparticle multipletts

 $G | \mathbf{f} \overline{\mathbf{f}} \rangle = (-1)^{I} C | \mathbf{f} \overline{\mathbf{f}} \rangle = (-1)^{I+L+S} | \mathbf{f} \overline{\mathbf{f}} \rangle$ $G | \pi^{\pm 0} \rangle = (-1)^{1} C | \pi^{\pm 0} \rangle = - | \pi^{\pm 0} \rangle$ $G | \rho \rangle = (-1)^{1} C | \rho \rangle = + | \rho \rangle$ $G | \omega \rangle = (-1)^{0} C | \omega \rangle = - | \omega \rangle$ $G | \sigma \rangle = (-1)^{0} C | \sigma \rangle = + | \sigma \rangle$

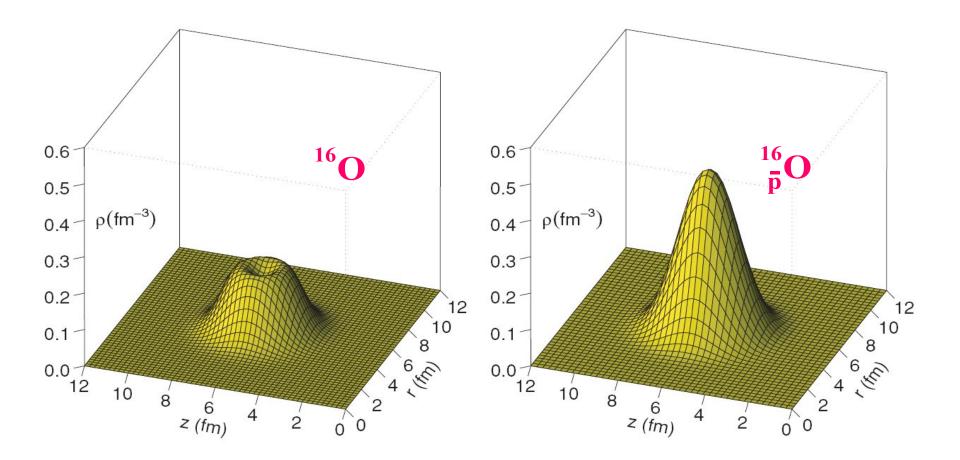
- Hans-Peter Dürr and Edward Teller, Phys. Rev. 101, 494 (1956)
 - sign change in coupling constant

$$V(NN)(r) = \sum_{M} V_{M}(r) \rightarrow V(N\overline{N})(r) = \sum_{M} G_{M}V_{M}(r)$$



Why do antihyperons in matter matter

- ► strong cold compression ⇒ color degrees of freedom might become very important
 - ▶ I.N. Mishustin *et al.*, PRC 71, 035201 (2005)



Elastic Antiproton-Nucleus Scattering

Elastic Scattering of Antiprotons from Complex Nuclei*

Gerson Goldhabert and Jack Sandweiss‡

Physics Department and Radiation Laboratory, University of California, Berkeley, California (Received May 5, 1958) pantiwith $\geq 2^{\circ}$.)

TABLE III. Comparison of experimental data for elastic antiproton-nucleus scattering of energy $T_{\bar{p}}=80$ to 200 Mev with Glassgold's calculations at $T_{\bar{p}}=140$ Mev. (Projected angle $\geq 2^{\circ}$.)

Angular interval (degrees)	Experimental $(T_{\overline{p}} = 80 \text{ to} 200 \text{ Mev})$	Number of events Calculated for V = -15 Mev W = -50 Mev	S Calculated for V = -528 MeV W = -50 MeV		
$\begin{array}{r} 2-6\\ 6-12\\ 12-24\\ 24-180\end{array}$	54 20 5 1	$56 \\ 17.1 \\ 4.3 \\ 1.4$	71 24 10 9.5		
2–180	80	78.8	114.5		

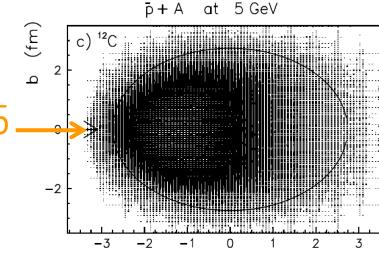
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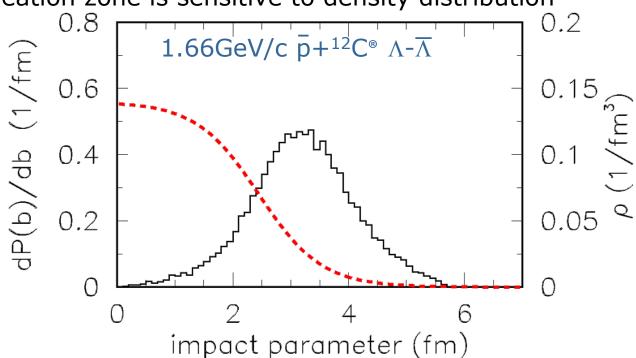
Annihilation Zone

- Anti-baryons are strongly absorbed:
 - ▶ p: 50 mb
 - Λ: 20 mb
 - $\overline{\Lambda}$: 100 mb/(1+p_{\overline{\Lambda}}/GeV)
- Emission of AA pair is anisotropic in center-of-mass
- Coincident detection of Λ and Λ constraints annihilation points



Sibirtsev et al

Creation zone is sensitive to density distribution



The potentials

- Fermi momentum p_F=220MeV/c
- Λ , $\overline{\Lambda}$ angular distribution of free events $\overline{p} + p^{\circ} \Lambda + \overline{\Lambda}$
- Potentials and propagation

$$(E - V)^{2} = (M_{0} + S)^{2} + \vec{P}_{in}^{2}$$
$$\vec{P}_{out}^{2} + M_{0}^{2} = \left(\sqrt{(M_{0} + S)^{2} + \vec{P}_{in}^{2}} + V\right)^{2}$$

Momentum independent

Potential at ρ_0	р	p	Λ
V [MeV]	300	200	200
S [MeV]	-342	-342	-228
V+S [MeV]	-42	-142	-28

- Potentials scaled linearly with local nucleon density
- \blacktriangleright No compression by the presence of the Y and \overline{Y}
- Potentials act instantaneously at the \bar{p} annihilation time



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Choice of Potentials

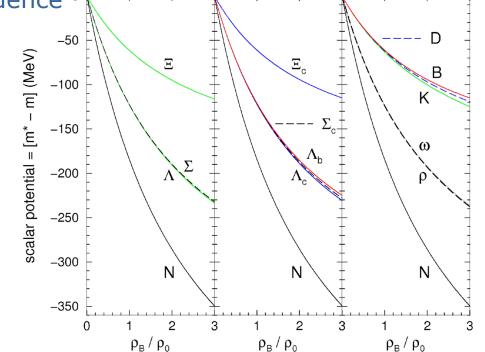


- ▶ K. Saito, K. Tsushima, A.W. Thomas, Prog. Part. Nucl. Phys. 58 (2007) 1
- - Scalar Potential

$$U_{V} = U_{V}(\rho_{0}) \cdot \frac{\rho}{\rho_{0}}$$

• Vector Potentia
• N=Z:

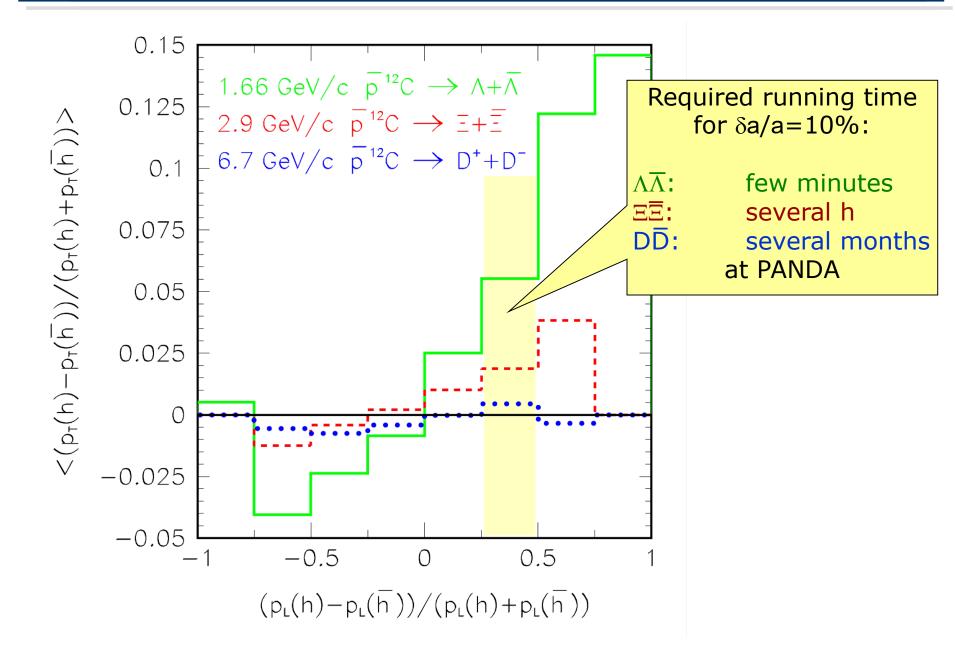
$$U_{V} = 41.8 \cdot (n_{q} - n_{\overline{q}}) \frac{\rho}{\rho_{0}}$$



potential	_	_		_	_						
V [MeV]	300	200	200	125	-125	84	-84	42	-42	-42	42
S [MeV]	-342	-342	-228	-184	-184	-123	-123	-61	-61	-61	-61
V+S [MeV]	-42	-142	-28	-59	-309	-39	-207	-19	-103	-103	-19

Other hadron-antihadron pairs

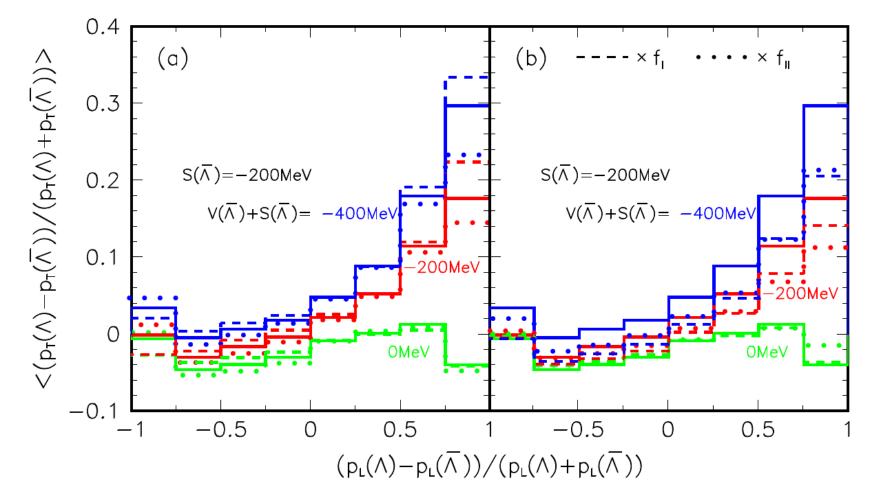




Momentum dependent potential

- Potentials may be momentum dependent
- ► Scale potential according to Gale or Lee and co-workers

$$f_I = \frac{1}{1 + (p - p_F)^2 / \Lambda^2}$$
 $f_{II} = \frac{1}{1 + \alpha \cdot (p/p_F)^{\beta}}$



Some open questions



rescattering

- influence of nuclear mass \Rightarrow use light nucleus to reduce rescattering
- ▶ but: coherence length of Λ anti Λ pair: t~ \hbar/E_F ~5fm/c \Rightarrow need large nucleus
- use Λ and anti-Λ polarization to enhance anti-ΛΛ pairs which did not encounter a rescattering on their way out

Summary

Antiproton collisions with nuclei are the ideal tool to produce exclusive hyperon-antihyperon pairs in nuclei at moderate momenta

Transverse momentum correlations of hyperon-antihyperon pairs produced close to threshold offer a unique opportunity to explore the potential of antihyperons relative to that of hyperons

Existing data of PS185 are encouraging e.g. fot PANDA, but not conclusive

Many improvements needed/open questions

- Momentum dependence of potentials
 - ▷ Reduce effect particularly for $\Xi \overline{\Xi}$
 - \triangleright Angular dependence of $\alpha_T \Leftrightarrow$ study momentum-dependence

Rescattering

- \triangleright influence of nuclear mass \Rightarrow use light nucleus to reduce rescattering
- Formation time
 - ▷ coherence length of $\Lambda \overline{\Lambda}$ pair: t~ $\hbar/E_F \sim 5 \text{fm/c} \Leftrightarrow \text{nuclei}$ of different size

5 decades of hyperons in neutron stars

NEUTRON STAR MODELS

A. G. W. CAMERON Atomic Energy of Canada Limited, Chalk River, Ontario, Canada Received June 17, 1959

Another reason why the writer has not taken into account complications inherent in using a relativistic equation of state is that no such things as pure neutron stars can be expected to exist. The neutrons must always be contaminated with some protons and sometimes with other kinds of nucleons (hyperons or heavy mesons).

Alastair G.W. Cameron, Astrophysical Journal, vol. 130, p.884 (1959)

Hyperons in neutron stars (2008)



Haris Djapo, Bern-Jochen Schäfer and Jochen Wambach arXiv:0811.2939v1 [nucl-th] 18 Nov 2008

In conclusion, irrespective of the YN interactions, incompressibility and symmetry parameter used, hyperons will appear in dense nuclear matter at densities around $\sim 2\rho_0$. This immediately leads to a softening of the EoS which in turn results in a smaller maximum mass of a neutron star.

With the prediction of a low onset of hyperon appearance it becomes practically impossible to ignore strangeness when considering neutron stars. Even though the prediction for the maximum masses of neutron stars are too low, the appearance of hyperons in neutron stars is necessary and any approach to dense matter must address this issue.

Baryon stars

10[°]

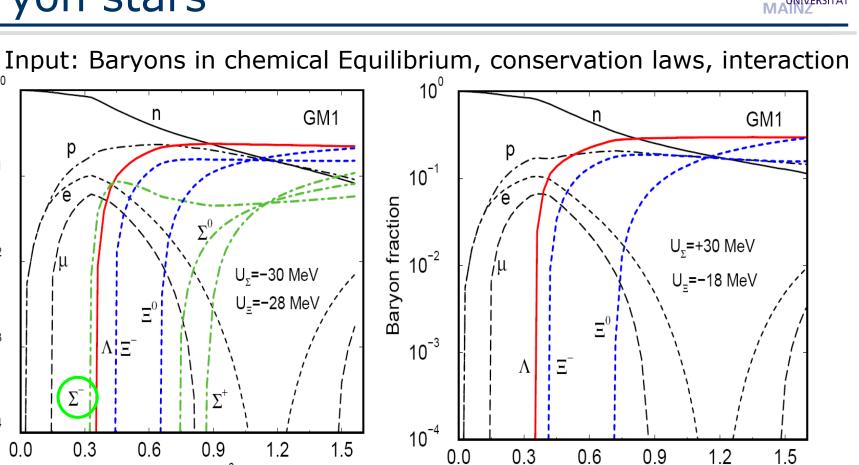
10⁻¹

10⁻²

10⁻³

10⁻⁴

Baryon fraction



N. K. Glendenning, Phys. Rev. C 64, 025801 (2001)

Density (fm^{-3})

- beyond $2\rho_0$ hyperons may play a significant role in neutron stars
- in the core hyperons may even be more abundant than neutrons
- needed: BB interaction at high density = at small distances

Density (fm^{-3})